

# The Power Series Method for Nonlocal and Nonlinear Evolution Equations

# Dr.Madhaba Chandra Rout

Department of Basic Science, Aryan Institute of Engineering & Technology, Bhubaneswar

# Dr. C K Sharma

Department of Basic Science, NM Institute of Engineering & Technology, Bhubaneswar

### Itishree Swain

Department of Basic Science, Raajdhani Engineering College, Bhubaneswar

ABSTRACT. The initial value problem for a 4-parameter family of nonlocal and nonlinear evo-lution equations with data in a space of analytic functions is solved by using a power series method in abstract Banach spaces. In addition to determining the power series expansion of the solution, this method also provides an estimate of the analytic lifespan expressed in terms of the norm of the initial data, thus establishing an abstract Cauchy-Kovalevsky type theorem for these equations.

#### I. INTRODUCTION AND RESULTS

In this work we prove an abstract Cauchy-Kovalevsky theorem for the following 4-parameter family of Camassa-Holm type equations

$$u + u^{k}u - au^{k-2}u^{3} + Q (1 - Q^{2})^{-1} \frac{b}{k+1} u^{k+1} + c u^{k-1}u^{2}_{x} - a(k-2) u^{k-3}u^{4}_{x}$$

$$+ (1 - \partial^{2})^{-1} k(k+2) - 8a - b - c(k+1) u^{k-2}u^{3}_{x} - 3a(k-2) u^{k-3}u^{3}_{x}u_{xx} = 0, \quad (1.1)$$

which was introduced in [HMa1] and is referred there as the k-abc-equation. The three parameters a, b and c range over the real numbers while k is a positive integer, whose value depends on a. If  $a \neq 0$  the  $b \geq 2$  and the presence of the term  $au^{k-2}u_x^3$  makes k-abc-equation a nonlocal and nonlinear equation which is not quasilinear. For k=2 and k=20 and k=21, we obtain the abfamily of equations (ab-equation) with cubic nonlinearities

$$u_t + u^2 u_x - a u_x^3 + \partial_x \left(1 - \partial_x^2\right)^{-1} \frac{b}{3} u^3 + \frac{6 - 6a - b}{2} u u_x^2 + \left(1 - \partial_x^2\right)^{-1} \frac{2a + b - 2}{2} u_x^3 = 0, \quad (1.2)$$

which was also introduced in [HMa1] and which contains two well-known integrable equations with cubic nonlinearities. In fact, for a = 1/3 and b = 2 the ab-equation gives the Fokas-Olver-Rosenau-Qiao (FORQ) equation (also known as the modified Camassa-Holm equation)

$$\partial_t u + u^2 \partial_x u - \frac{1}{3} (\partial_x u)^3 + \partial_x (1 - \partial_x^2)^{-1} \frac{2}{3} u^3 + u (\partial_x u)^2 + (1 - \partial_x^2)^{-1} \frac{1}{3} (\partial_x u)^3 = 0, \quad (1.3)$$

which was derived in different ways by Fokas [F], Olver and Rosenau [OR] and Qiao [Q], and also appeared in a work by Fuchssteiner [Fu]. For a = 0 and b = 3 the ab-equation gives the

Novikov equation (NE)

$$u_t + u^2 u_x + \partial_x (1 - \partial_x^2)^{-1} u^3 + \frac{3}{2} u u_x^2 + (1 - \partial_x^2)^{-1} \frac{1}{2} u_x^3 = 0,$$
 (1.4)

which was derived by V. Novikov in [N1], where he provides a classification of all integrable CH-type equations with quadratic and cubic nonlinearities.

Finally, for a=0 and c=(3k-b)/2 the k-abc-equation makes sense for all  $k\geq 1$  and gives the following generalized Camassa-Holm equation (g-kbCH)

$$u + u^{k}u + (1 - \partial^{2})^{-1}\partial \frac{b}{x} u^{k+1} + \frac{3k - b}{2} u^{k-1}u^{2} + (1 - \partial^{2})^{-1} \frac{(k-1)(b-k)}{x} u^{k-2}u^{3} = 0,$$

$$t + u^{k}u + (1 - \partial^{2})^{-1}\partial \frac{b}{x} u^{k+1} + \frac{3k - b}{2} u^{k-1}u^{2} + (1 - \partial^{2})^{-1} \frac{(k-1)(b-k)}{x} u^{k-2}u^{3} = 0,$$

$$(1.5)$$

www.ijceronline.com Open Access Journal Page 64

Date: March 17, 2016. \*Corresponding author: himonas.1@nd.edu.

<sup>2010</sup> Mathematics Subject Classification. Primary: 35Q53, 37K10, 35C07.

Key words and phrases. Abstract Cauchy-Kovalevsky and Ovsyannikov theorems, nonlinear and nonlocal evolution equations, analytic spaces and lifespan, integrable equations, Camassa-Holm type equations, peakon traveling wave solutions.

which is a quasilinear equation with (k + 1) order nonlinearities and which was studied in [HH2] and [GH]. When k = 1 the g-kbCH equation gives the well-known b-equation

$$\underbrace{m}_{X} + \underbrace{um}_{X} + \underbrace{bu}_{X} \underbrace{m}_{X} = 0, \quad m = u - u_{xx}$$
evolution convection stretching (1.6)

having quadratic nonlinearities. In this local form it was introduced by Holm and Staley [HS1, HS2] and it expresses a balance between evolution, convection and stretching. The b-equation (1.6) contains two integrable members, namely the Camassa-Holm (CH) equation that corresponds to b=2 and the Degasperis-Procesi (DP) equation that corresponds to b=3. Mikhailov and Novikov [MN] proved that there are no other integrable members of the b-equation. Furthermore, V. Novikov [N2] recently proved that the only other integrable member of the g-kbCH equation (1.5) apart from CH and DP is the NE (1.4). To summarize, thus far it is known that the k-abc-equation (1.1) contains four integrable equations, CH, DP, FORQ and NE. However, the existence of other integrable members of the k-abc-equation remains an open question. In fact, integrability theory provides one of the motivations for studying such nonlocal equations.

Another motivation for studying equations like *k-abc*-equation is the quest for equations capturing wave breaking and peaking, which goes back to Whitham, who articulates it in his 1974 book [W] (p. 477) as follows: "Although both breaking and peaking, as well as criteria for the occurrence of each, are without doubt contained in the equations of the exact potential theory, it is intriguing to know what kind of simpler mathematical equation could include all these phenomena." It is remarkable that the *k-abc*-equation has peakon traveling wave solutions for all values of the four parameters *k*, *a*, *b* and *c*. These, including multipeakons, have been derived in [HMa1]. The peakon solutions in the non-periodic case can be written in the following form

$$u(x, t) = ye^{-|x-(1-a)y^k t|}, y \in \mathbb{R}.$$
 (1.7)

When a=0 and c=(3k-b)/2, which is the case of the generalized Camassa-Holm equation (1.5), these are of the form  $u(x,t)=c^{1/k}e^{-|x-ct|}$  and were derived in [GH], together with the corresponding multipeakon on the line and the circle. In the case k=1, which is the b-equation, peakon solutions were derived by Holm and Staley [HS1, HS2], who made the important observation that the b-equation has peakon (and multipeakon) traveling wave solutions for all values of b. Of course, it is Camassa and Holm [CH] who observed first that the celebrated CH equation has the peakon (weak) solutions  $u(x,t)=ce^{-|x-ct|}$ .

The initial value problem of the k-abc-equation in Sobolev spaces was studied in [HMa2]. More precisely, there the following well-posedness result was obtained. If  $a,b,c\in\mathbb{R}$  with  $a\neq 0$  and  $k\in\mathbb{N}$  with  $k\geq 2$ , then the Cauchy problem for the k-abc-equation (1.1) with initial data  $u(x,0)=u_0(x)\in H^s$ ,  $x\in\mathbb{R}$  or T,  $s>\frac{5}{\sqrt{2}}$  has a unique solution  $u\in C([0,T];H^s)$ . Furthermore, the lifespan T satisfies the estimate

$$T = \frac{1}{\|u_0\|_{H^s}^k}.$$
 (1.8)

Also, in [HMa2] continuity properties of the data-to-solution map are investigated. The case a=0 and c=(6-6a-b)/2 yields the g-kbCH equation, whose well-posedness in Sobolev spaces  $H^s$  for all  $s>\frac{3}{2}$ was proved in [HH2]. The Cauchy problem for the integrable members CH, DP, FORQ and NE was studied earlier by many authors. For some results on well-posedness, traveling wave solutions, and other analytic properties of these and related equations we refer the reader to the following works and the references therein [BHP2], [BSS], [BC], [CHT], [CK], [CL], [CM], [CS], [DP], [EY], [FF], [HH1], [HK], [HMPZ], [KL], [LO], [LS], [R], [B], [CKSTT], [KPV], [GLOQ], [HHG], [HGH], [HLS], [KT], [L], and [Ti].

In this work we study the initial value problem for k-abc-equation when the initial data belong to spaces of analytic functions, on both the line and the circle. More precisely, these spaces are defined as follows. For  $s \ge 0$  and  $\delta > 0$ , on the line these spaces are defined by

defined as follows. For 
$$s \ge 0$$
 and  $\delta > 0$ , on the line these spaces are defined by 
$$G^{\delta,s}(\mathbb{R}) = \{ \phi \in L^2(\mathbb{R}) : ||\phi||_{G^{\delta,s}(\mathbb{R})}^2 \doteq \int_{\mathbb{R}} \langle \xi \rangle^{2s} e^{2\delta|\xi|} |\phi(\xi)|^2 d\xi < \infty \}, \tag{1.9}$$

www.ijceronline.com Open Access Journal Page 65

where  $\langle \xi \rangle$   $(1+\xi^2)^{1/2}$ . On the circle the corresponding spaces are  $\geq$ 

$$G^{\delta,s}(T) = \{ \phi \in L^{2}(T) : ||\phi||_{G^{\delta,s}(T)}^{2} \stackrel{\wedge}{=} \langle k \rangle^{2s} e^{2\delta |k|} |\phi(k)|^{2} < \infty \}, \tag{1.10}$$

where  $\langle k \rangle \doteq (1+k^2)^{1/2}$ . Here, when a result holds for both on the line and the circle then we use the notation  $||\cdot||_{\delta,s}$  for the norm and  $G^{\delta,s}$  for the space in both cases. We observe that a function  $\phi$  in  $G^{\delta,s}$  has an analytic extension to a symmetric strip around the real axis with width  $\delta$  (see Lemma 2). This  $\delta$  is called the *radius of analyticity* of  $\phi$ . Since in the next Theorem we will assume that the initial data  $u_0$  is in  $G^{1,s+2}$  we would like to point out that in the periodic case an analytic function (i.e. an element of  $C^\omega(T)$ ) belongs to a  $G^{\delta_0,s}(T)$ , for some  $\delta_0 > 0$  and any  $s \ge 0$ . More precisely we have the following result, whose proof is easy and will be omitted.

**Lemma 1.** If  $u_0 \in C^{\omega}(T)$ , there exists  $\delta_0 > 0$  such that  $u_0 \in G^{\delta_0,s}(T)$  for any  $s \ge 0$ .

Next, we state our main result, which is motivated by [BHP1] and [BHP2]. For the sake of simplicity we shall assume that our initial data  $u_0$  belong in  $G^{1,s+2}$ .

**Theorem 1.** Let  $s > {}^1._{\overline{z}} f u_0 \in G^{1,s+2}$  on the circle or the line, then there exists a positive time T, which depends on the initial data  $u_0$  and s, such that for every  $\delta \in (0,1)$ , the Cauchy problem for the k-abc-equation (1.1) with initial condition  $u(x,0) = u_0(x)$  has a unique solution  $u(x,0) = u_0(x)$  and  $u(x,0) = u_0(x)$  has a unique solution  $u(x,0) = u_0(x)$  and  $u(x,0) = u_0(x)$  has a unique solution  $u(x,0) = u_0(x)$  and  $u(x,0) = u_0(x)$  has a unique solution  $u(x,0) = u_0(x)$  and  $u(x,0) = u_0(x)$  has a unique solution  $u(x,0) = u_0(x)$  has a unique solution  $u(x,0) = u_0(x)$  and  $u(x,0) = u_0(x)$  has a unique solution  $u(x,0) = u_0(x)$  ha

$$T = \frac{1}{||u_0||_{1,s+2}^k}. (1.11)$$

A more precise statement of estimate (1.11) is provided in Section 4 (see (4.13)). For the Camassa-Holm equation on the circle, a result similar to Theorem 1 but without an analytic lifespan estimate like (1.11) was proved in [HM]. Furthermore, for CH, DP, NE and FORQ Theorem 1 was proved in [BHP2] using a different approach based on a contraction type argument in an appropriate space which is build from a scale of Banach spaces. Here we are using a power series method, which for the quasilinear g-kbCH equation was presented in [BHP1]. The novelty of this work is that it provides a comprehensive treatment of a large family of Camassa-Holm type equations whose local part includes non-quasilinear terms like  $u^{k-2} q^3$ . Furthermore, it makes a complete presentation of the autonomous Ovsyannikov theorem using the power series method, on which the proof of Theorem 1 is based. We conclude, by mentioning that there are many versions of the abstract Cauchy-Kovalevsky theorem proved in a variety of ways. Many of these works are motivated by water wave models and the Euler equations. For more information about these we refer the reader to Baouendi and Goulaouic [BG1], [BG2], Ovsyannikov [O1], [O2], Treves [Tre1], [Tre2], Nirenberg [Nr], Nishida [Ns], Caflisch [C], Safonov [S], and the references therein.

The paper is organized as follows. In Section 2 we describe the spaces used and prove the needed properties. In Section 3, following the work of Treves, we presents the proof of the autonomous Ovsyannikov theorem which solves an abstract Cauchy problem by the power series method. Finally, in Section 4 we apply the power series method to the *k-abc*-equation and derive the estimates needed for this method.

2. 
$$G^{\delta,s}$$
 spaces and results

We begin with the properties of the  $G^{\delta,s}$  and the estimates needed to prove our main result. The next two lemmas provide a better understanding of the spaces  $G^{\delta,s}$  and their properties. One can easily prove these results.

**Lemma 2.** Let  $\phi \in G^{\delta,s}$ . Then,  $\phi$  has an analytic extension to a symmetric strip around the real axis of width  $\delta$ , for  $s \ge 0$  in the periodic case and  $s > \frac{1}{2}$  in the non-periodic case.

**Lemma 3.** If  $0 < \delta^{s} < \delta \le 1$ ,  $s \ge 0$  and  $\phi \in G^{\delta,s}$  on the circle or the line, then

$$||\partial_x \phi||_{\delta', s} \le \frac{e^{-1}}{\delta - \delta^{\dagger}} ||\phi||_{\delta, s} \tag{2.1}$$

$$||\partial_x \phi||_{\delta,s} \le ||\phi||_{\delta,s+1} \tag{2.2}$$

$$||(1 - \partial_{x}^{2})^{-1} \phi||_{\delta, s+2} = ||\phi||_{\delta, s}$$
(2.3)

$$||(1-\partial_x^2)^{-1}\phi||_{\delta,s} \le ||\phi||_{\delta,s}$$
 (2.4)

$$||\partial_x (1 - \partial_y^2)^{-1} \phi||_{\delta,s} \le ||\phi||_{\delta,s}. \tag{2.5}$$

Furthermore, we shall need to prove an algebra property for these spaces, which is the main result in the following lemma.

**Lemma 4.** For  $\phi \in G^{\delta,s}$  on the circle or the line the following properties hold true:

1) If 
$$0 < \delta^{|} < \delta$$
 and  $s \ge 0$ , then  $||\cdot||_{\delta,s}^2 \le ||\cdot||_{\delta,s}^2$ ; i.e.  $G^{\delta,s} \longleftrightarrow G^{\delta',s}$ .

2) If 
$$0 < s^{1} < s$$
 and  $\delta > 0$ , then  $||\cdot||_{\delta,s}^{2} \le ||\cdot||_{\delta,s}^{2}$ ; i.e.  $G^{\delta,s} \longleftrightarrow G^{\delta,s'}$ .

3) For s > 1/2 and  $\phi, \psi \in G^{\delta,s}$  we have

 $\frac{||\phi\psi||_{\delta,s} \leq c_s ||\phi||_{\delta,s}||\psi||_{\delta,s}}{\|\phi\psi\|_{\delta,s} \leq c_s ||\phi||_{\delta,s}||\psi||_{\delta,s}} \qquad (2.6)$ where  $c_s = 2(1+2^{2s})^{\sum_{\infty} \frac{1}{\langle \xi \rangle}} \text{ in the periodic case and } c_s = 2(1+2^{2s})^{\int_{0}^{\infty} \frac{1}{\langle \xi \rangle}} d\xi \text{ in the periodic case}$ non-periodic case.

Proof of Lemma 4. We will provide the proof in the periodic case. The proof in the nonperiodic case is similar. Properties (1) and (2) follow directly from the definition of the spaces  $G^{\delta,s}$  and the corresponding norms. Therefore, we restrict our attention to the proof of the algebra property (3), which reads as follows

$$||\phi\psi||_{\delta,s}^{2} = \sum_{k \in \mathbb{Z}} \langle k \rangle^{2s} e^{2\delta|k|} ||\phi\psi(k)||^{2} = ||\langle k \rangle^{s} e^{\delta|k|} ||\phi\psi(k)||_{\mathbb{Z}_{2}}^{2}$$

$$= ||\langle k \rangle^{s} e^{\delta|k|} \sum_{k \in \mathbb{Z}_{2}} ||\phi(n)\psi(k-n)||_{\mathbb{L}_{2}(\mathbb{Z}_{2})}^{2} \leq c_{s} ||\phi||_{\delta,s} ||\psi||_{\delta,s}. \tag{2.7}$$

Defining f and g by  $f(k) = \langle k \rangle {}^s e^{\delta_j k} | \phi(k) \rangle {}^n$  and  $g(k) = \langle k \rangle {}^s e^{\delta_j k} | \psi(k) \rangle {}^n$  when see that the algebra property (2.7) is equivalent to

$$\stackrel{...}{\underset{e}{\sim}} k)^{s} \stackrel{\delta|k|}{\underset{n \in \mathbb{Z}}{\sim}} \frac{e^{-\delta|n|}f(n)}{\binom{s}{s}} \frac{e^{-\delta|k-n|}g(k-n)}{\binom{s}{s}} \frac{c^{2}}{\underset{n \in \mathbb{Z}}{\sim}} \frac{c^{2}}{\underset{n \in \mathbb{Z}}{\sim}} ||f||_{\mathcal{U}(\mathbb{Z})}^{2} ||g||_{\mathcal{U}(\mathbb{Z})}^{2}.$$
 (2.8)

Furthermore, using the triangle inequality  $|k| \le |n| + |k - n|$  we notice that in order to prove (2.8) it suffices to show that

Using the Cauchy-Schwarz inequality we have

$$\sum_{n\in\mathbb{Z}}\frac{f(n)}{\langle n\rangle}\frac{g(k-n)}{s\langle k-}^2 \leq \sum_{n\in\mathbb{Z}}\frac{1}{\langle n\rangle^{2s}\langle k-} \sum_{n\in\mathbb{Z}}|f(n)|^2|g(k-n)|^2,$$

which gives the following bound for the left-hand side of (2.9

$$\leq \sup_{k \in \mathbb{Z}} \left\langle k \right\rangle^{2S} \frac{1}{\langle n \rangle^{2S} \langle k - \rangle} \sum_{n \in \mathbb{Z}} \left\{ f(n)^{2} g(k - n)^{2} \right\}^{1/2} \frac{1}{\langle n \rangle^{2S} \langle k - \rangle} \left( 2.10 \right)$$

Now, we need the following estimate, whose proof is given after the one of algebra property.

**Lemma 5.** For 
$$l > 1/2$$
,  $a \in \mathbb{Z}$  and  $c_l = 2(1 + 2^{2l}) \sum_{k=0}^{\infty} \frac{1}{\langle a \rangle^{2l}}$  we have 
$$\sum_{k \in \mathbb{Z}} \frac{1}{\langle k \rangle^{2l} \langle k - a \rangle^{2l}} \leq \langle a \rangle^{2l}. \tag{2.11}$$

In the non-periodic case, for 
$$a \in \mathbb{R}$$
 and  $c_{f} = \frac{1}{2(1+2^{2f})^{\int_{\xi \geq 0} \frac{1}{-1} d\xi} d\xi}$  we have
$$\int \frac{1}{\mathbb{R} \langle x \rangle^{2f} \langle x - a \rangle^{2f}} dx \leq \frac{c_{f}^{2}}{\langle a}. \tag{2.12}$$

Since s > 1/2, combining (2.10) and (2.11) and interchanging the order of summations (Fubini's theorem) we get

$$\begin{aligned} & ...\langle k \rangle^{s} \underbrace{\sum_{n \in \mathbb{Z}} f(n) g(k - \frac{1}{s}) g(k - \frac{1}{s})}_{n \setminus s} g(k - \frac{1}{s}) g($$

which completes the proof of Lemma 4. Q

**Proof of Lemma 5.** If a = 0 then it is easy to see that inequality (2.11) holds true. If a is a negative integer then making the change of variables m = k - a the estimate is reduced to the case where a is replaced with -a. Thus, it suffices prove inequality (2.11) only when a is a positive integer. For this we shall use the notation [x], which stands for the biggest integer that is less than or equal to x. Thus,  $x - 1 < [x] \le x$ . Now, we begin by decomposing our sum as

$$S \stackrel{:}{=} \underbrace{\sum_{k \in \mathbb{Z}} \frac{1}{\langle k \rangle^{2l} \langle k - |}}_{k \in \mathbb{Z}} = \underbrace{\sum_{-\infty < k \le 0} \frac{1}{\langle k \rangle^{2l} \langle k - a \rangle^{2l}}}_{-\infty < k \le 0} \underbrace{\sum_{-\infty < k \le 0} \frac{1}{\langle k \rangle^{2l} \langle k - a \rangle^{2l}}}_{S_1 \le k \le [a/2]} + \underbrace{\sum_{1 \le k \le [a/2]} \frac{1}{\langle k \rangle^{2l} \langle k - a \rangle^{2l}}}_{[a/2] + 1 \le k \le [a]} \underbrace{\sum_{1 \le k \le [a]} \frac{1}{\langle k \rangle^{2l} \langle k - a \rangle^{2l}}}_{S_1 + S_2 + S_3 + S_4.}$$

$$Sobtain$$

By setting  $\vartheta = -k$  we obtain

$$S_{1} = \sum_{-\infty < k \le 0} \frac{1}{(1+k^{2})^{l}(1+(k-a)^{2})^{l}} \le \frac{1}{\langle a \rangle} \sum_{\theta=0}^{\infty} \frac{1}{\langle \theta \rangle}.$$
By noticing that in  $S_{2}$  we have  $1+(k-a)^{2} \ge 1+a^{2}/4 = \frac{4+b^{2}}{4} > \frac{1+a^{2}}{4}$  we obtain

$$S_2 \leq \frac{2^{l}}{\langle a \rangle} \sum_{k=1}^{\infty} \frac{1}{\langle k \rangle}$$

We also notice that in  $S_3$ 

icing that in 
$$S_2$$
 we have  $1 + (k - a)^2 \ge 1 + a^2/4 = \frac{4 + b^2}{4} > \frac{1 + a^2}{4}$  we obtain
$$S_2 \le \frac{2}{\sqrt{2}} \frac{\sum_{j=1}^{\infty} 1}{\sum_{j=1}^{2^{j}} \frac{\sum_{j=1}^{\infty} 1}{\sum_{j=1}^{2^{j}} \frac{\sum_{j=1}^{\infty} 1}{\sum_{j=1}^{2^{j}} \frac{\sum_{j=1}^{\infty} 1}{\sum_{j=1}^{\infty} \frac{\sum_{j=1}^{2^{j}} \frac{\sum_{j=1}^{\infty} 1}{\sum_{j=1}^{\infty} 1}}}}}}}$$

$$S_4 \leq \frac{1}{\langle a \atop 2l \atop m=1} \frac{\sum_{\substack{1 \\ 2l}} 1}{\langle m \rangle}.$$

Adding the estimates obtained above for the sums  $S_1$ ,  $S_2$ ,  $S_3$  and  $S_4$  gives inequality (2.11). In order to prove (2.12) one has just to replace the sums with integrals. This completes the proof of Lemma 5. Q

#### 3. THE POWER SERIES METHOD

Here, following Treves [Tre1], [Tre2], we prove the autonomous Ovsyannikov theorem using the power series method, which consists of finding a solution for the Cauchy problem

$$\frac{du}{dt} = F(u), \quad u(0) = u_0, \tag{3.1}$$

given by a formal power series

$$u(t) = \sum_{m=0}^{\infty} u_m t^m, \tag{3.2}$$

Open Access Journal www.ijceronline.com Page 68 and estimating the coefficients  $u_m$  in order to prove the convergence of the above series. We begin with the following important definition. Let  $\{X_\delta\}_{0<\delta\leq 1}$  be a decreasing scale of Banach spaces and  $u_0 \in X_1$  be given. Also, let  $X_0 = X_{\delta}$ .

**Definition 1.** We say that  $F: X_0 \to X_0$  is Ovsyannikov analytic at  $u_0$  if there exist positive constants R, A and  $C_0$  such that for all  $k \in \mathbb{Z}_+$  and  $0 < \delta^{\scriptscriptstyle \parallel} < \delta < 1$  we have

$$||D^{k}F(u)(v_{1},...,v_{k})||_{\delta'} \leq \frac{A}{\delta} \frac{C^{k}k!}{\delta - \frac{0}{\delta^{1}}} ||v_{1}||_{\delta} ... ||v_{k}||_{\delta'}$$
(3.3)

 $||D^{k}F(u)(v_{1},...,v_{k})||_{\delta'} \leq \frac{A}{\delta} \frac{C^{k}k!}{\delta - \frac{0}{\delta!}} ||v_{1}||_{\delta} ... ||v_{k}||_{\delta'}$ for all  $u \in \{u \in X_{\delta} : ||u - u_{0}||_{\delta} < R\}$  and  $(v_{1},...,v_{k}) \in X_{\delta}^{k} = \underbrace{X_{\delta} \times \cdots \times X_{\delta}}_{L'}$ , where  $D^{k}F$  is the

Frechet derivative of F of order k.

Here, we shall prove the following important result.

**Theorem 2.** If  $u_0 \in X_1$  and F is Ovsyannikov analytic as above, then there exists T > 0 such that the Cauchy problem (3.1) has a unique solution which, for every  $\delta \in (0, 1)$  is a holomorphic function in  $D(0,T(1-\delta))$  valued in  $X_{\delta}$  satisfying

$$\sup_{|t| < T(1-\delta)} \|u(t) - u_0\|_{\delta} < R, \quad 0 < \delta < 1.$$
 (3.4)

Moreover, the lifespan T is given by

$$T = \frac{1}{2e^2AC_0},$$

where the constants R, A and  $C_0$  come from the Definition 1.

From the definition of Ovsyannikov analyticity one can easily prove the following result.

**Proposition 1.** If F is Ovsyannikov analytic at  $u_0 \in X_1$  then there is R > 0 such that, given any pair  $(\delta, \delta^i)$ ,  $0 < \delta^i < \delta < 1$  and any  $u \in B_{\delta}(u_0; R)$  the Taylor series

$$\sum_{k=0}^{\infty} \frac{1}{k} p^k F(u^0) \left( \underline{u - u^0, \dots, u - u^0} \right)$$

converges absolutely to F(u) in  $X_{\delta'}$ .

**Proof of Theorem 2.** Let  $0 < \delta^1 < \delta < 1$ . Since  $u_0 \in X_1$  and F is Ovsyannikov it follows from Proposition 1 that

$$F(u) = \sum_{k=0}^{\infty} \frac{1}{k!} P^k F(u^0) \left( u - u^0, \dots, u - u^0 \right)$$
 in  $X_{\delta}$ .

Since we want to have  $\frac{du}{dt} = F(u)$ , we must have

$$u_{1} + \sum_{m=1}^{\infty} (m+1)u_{m+1}t^{m}$$

$$= F(u_{0}) + \sum_{m=1}^{\infty} \sum_{k=1}^{\infty} \sum_{m+1}^{\infty} \sum_{k=1}^{\infty} \frac{t^{m}}{n!} D^{k}F(u_{0})(u_{1},...,u_{m_{k}}).$$

$$k! \qquad 0 \qquad m_{1},...,u_{m_{k}}).$$

We then conclude that the coefficients of the series are given recursively by

$$u_1 = F(u_0) (3.5)$$

and

$$(m+1)u_{m+1} = \sum_{\substack{k=1 \ m_1 + \dots + m_k = m \\ m_l \ge 1}} \frac{D^k F(u_0)}{k!} (u_{m_1}, \dots, u_{m_k}), \ m \ge 1.$$
 (3.6)

We shall use the following lemma.

**Lemma 6.** For all k = 1, 2, ..., we have

$$\sum_{\substack{m_1 + \dots + m_k = m \\ m_l \ge 1}} \frac{1}{m_1^2 \dots m_k^2} \le \frac{1}{m^2} \frac{2\pi^2}{3}^{k-1}.$$

**Proof.** We prove the lemma by induction on k. For k = 1 the inequality is clear:

$$\sum_{\substack{m_1 + \dots + m_k = m \\ m_i \ge 1}} \frac{1}{m_1^2 \dots m_k^2} = \frac{1}{m^2} = \frac{1}{m^2} = \frac{2\pi^2}{3}^0.$$

For  $k \ge 2$  we can write

$$\sum_{\substack{m_1+\dots+m_k=m\\m_\ell\geq 1}} \frac{1}{m^2 \dots m^2} = \sum_{k=1}^{m-k+1} \frac{1}{k^2} \sum_{\substack{m_1+\dots+m_{k-1}=m-l\\m_\ell\geq 1}} \frac{1}{m^2 \dots m^2}...m^2.$$

Hence the induction hypothesis gives

$$\sum_{\substack{m_1+\dots+m_k=m\\ m_l\geq 1}} \frac{1}{m_1^2 \dots m_k^2} \leq \frac{2\pi^2}{3} \frac{k-2}{3} \frac{m_2^{-1}}{l^2(m-l)^2}.$$

We now see that

$$m^{2} \frac{\sqrt[m]{1}}{\sqrt{2(m-l)^{2}}} = \frac{m-k+1}{\sqrt{2(m-l)^{2}}} \frac{m^{2}}{\sqrt{2(m-l)^{2}}} = \frac{\sqrt[m]{2}}{\sqrt{2}} \frac{1}{\sqrt{l}} + \frac{1}{m-l}$$

$$\leq 2 \frac{\sqrt[m]{2}}{\sqrt{2}} \frac{1}{\sqrt{l}^{2}} + \frac{1}{(m-l)^{2}} \leq 4 \frac{2^{2}}{\sqrt{l}^{2}} \frac{1}{\sqrt{l}^{2}} = \frac{2\pi^{2}}{3}.$$

Therefore

$$\sum_{\substack{m_1 + \dots + m_k = m \\ m_l \ge 1}} \frac{1}{m_1^2 \dots m_k^2} \le \frac{2\pi^2}{3}^{k-1},$$

which proves the lemma. Q

We claim that

$$||u_m||_{\delta} \le \frac{\mu}{m^2} \frac{B}{1-\delta}^m$$
, for all  $m = 1, 2, ...$ , (3.7)

where  $\mu = \frac{3}{4C_0\pi^2}$  and  $B = 2e^2AC_0$ . To prove this fact, we proceed by induction on m. For m = 1, we know that  $u_1 = F(u_0)$ . Then it follows from (3.3) that

$$||u_1||_{\delta} = ||F(u_0)||_{\delta} \le \frac{A}{1-\delta} \le \frac{\mu B}{1-\delta}$$

which is true since  $\mu B = \frac{3e^2A}{2\pi^2} \ge A$ . For  $m \ge 2$ , we select  $v = \frac{1-\delta}{m+1} > 0$ . We notice that  $0 < \delta < \delta + \nu < 1$  and

$$\frac{1}{1 - (\delta + \nu)} = \frac{m+1}{m(1-\delta)} . \tag{3.8}$$

Thus, by using (3.6), the induction hypothesis and (3.8) we obtain

is, by using (3.6), the induction hypothesis and (3.8) we obtain 
$$\sum_{\Sigma} \sum \sum \sum AC_{0}^{k} ||u_{m_{1}}||_{\delta+\nu} \dots ||u_{m_{k}}||_{\delta+\nu}$$

$$\leq \frac{m+1}{1-\delta} \sum_{k=1 \ m_{1}+\dots+m_{k}=m} \frac{AC_{0}^{k}\mu^{k}}{m^{2}\dots m^{2}} \frac{B}{1-\delta} \sum_{m=1}^{m} 1 + \frac{1}{m}.$$

We then have

$$||u_{m+1}||_{\delta} \leq \frac{AC_0e\mu B^m}{m^2(1-\delta)^{m+1}} \sum_{k=1}^{\infty} \frac{2C_0\mu\pi^2}{3}^{k-1} = \frac{AC_0e\mu B^m}{m^2(1-\delta)^{m+1}} \sum_{k=1}^{\infty} \frac{1}{2}^{k-1},$$

since  $\mu = \frac{3}{4C_0\pi^2}$ . Hence,

$$||u_{m+1}||_{\delta} \leq \frac{2AC_0e\mu B^m}{m^2(1-\delta)^{m+1}} \leq \frac{2AC_0e^2\mu B^m}{(m+1)^2(1-\delta)^{m+1}} = \frac{\mu}{(m+1)^2} \frac{B}{1-\delta} \qquad ,$$

Open Access Journal Page 70 www.ijceronline.com

since  $B = 2AC_0e^2$ . The proof of the claim is complete.

Therefore, for  $T = \frac{1}{2e^2AC_0} = B^{-1}$  and for  $|t| \le T(1-\delta)$  the series  $u_m t^m$  converges absolutely

in  $X_\delta$  for the unique solution to our Cauchy problem. Q

#### 4. THE k-abc-EQUATION

In this section we apply the power series method for the *k-abc*-equation. For this we rewrite this equation in the following nonlocal form:

$$\partial_t u = F(u) = -(1 - \partial_x^2)^{-1} \partial_x \frac{b+1}{k+1} u^{k+1} + (c-k) u^{k-1} u_x^2 - u^k u_{xx} + 3a u^{k-2} u_x^2 u_{xx}$$

$$- (1 - \partial_x^2)^{-1} [k(k+2) - 9a - b - c(k+1)] u^{k-2} u^3 - {}_x 3a(k-2) u^{k-3} u^3 u_{xx}$$

$$(4.1)$$

We shall prove that there exist positive constants A, Co such that

$$||D^{j}F(u_{0})(v_{1},...,v_{j})||_{\delta',s+2} \leq \frac{AC^{j}j!}{\delta - \frac{0}{\delta^{j}}||v_{1}||_{\delta,s+2}...||v_{j}||_{\delta,s+2}},$$
(4.2)

for all  $(v_1,...,v_j) \in G^{\delta,s+2}$  and j=0,1,...,k+1 since for  $j \ge k+2$  we have

$$||D^{j}F(u_{0})(v_{1},\ldots,v_{i})||_{\delta',s+2}=0.$$

For simplicity, we shall provide estimate (4.2) only for the term

$$F_1(u) = (1 - \partial_v^2)^{-1} 3a(k - 2)u^{k-3}u^3u_{xx}$$

since the other terms can be estimated analogously. By using the following formula for the Frechet derivative of F of order j,  $1 \le j \le k+1$ , at the point  $u_0$ ,

$$D^{i}F(u_0)(v_1,\ldots,v_j)=\frac{d}{d\tau_j}\cdots\frac{d}{d\tau_1}\tilde{F}(u_0+\sum_{i=1}^{j-1}\tau_iv_i)\tilde{\tau}_{1=\cdots=\tau_j=0'}$$

we obtain, for  $v_l \in G^{\delta, s+2}$ , l = 1, ..., j,

$$D^{j} (u^{k-3} u_{x}^{3} u_{x}) \quad (v_{1}, \dots, v_{j}) =$$

$$= \frac{(k-3)!}{(k-3-j)!} u^{k-3-j} (\partial u)^{3} (\partial^{2} u) v v \dots v$$

$$(k-3-j)! \quad (4.3)$$

$$+\frac{(k-3)!}{(k-2-j)!}u^{k-2-j}(\partial u)^{3}\sum_{j=1}^{\infty}(\partial^{2}v)\tilde{v}$$

$$(4.4)$$

$$+3\frac{(k-3)!}{(k-2-j)!}u^{k-2-j}(\partial u)^{2}(\partial^{2}u) \sum_{i=1}^{\sum L} (\partial v)\tilde{v}$$

$$(4.5)$$

$$+3\frac{(k-3)!}{(k-1-j)!}u_{0}^{k-1-j}(\partial_{x}u)^{2}\sum_{i=1}^{j}\sum_{j=1}^{j}(\partial_{x}^{2}v)(\partial_{x}u)\tilde{v}_{i,j}$$
(4.6)

$$(k-2-j)! \xrightarrow{x \ 0 \ x \ 0} \xrightarrow{i=1} \xrightarrow{k-1} (k-3)! = \frac{(k-3)!}{(k-1-j)!} u^{k-1-j} (\partial u) (\partial^{2}u) \xrightarrow{j=1} \underbrace{\sum_{i=1}^{j} \sum_{j=1}^{j} \sum_{j=1}^{$$

$$+6\frac{(k-3)!}{(k-j)!}u^{k-j}(\partial u) \qquad \qquad (\partial^2 v)(\partial v)(\partial v)\tilde{v} \qquad (4.8)$$

$$(k-j)! \qquad \qquad (k-j)! \qquad \qquad (4.8)$$

$$+6\frac{(k-3)!}{u^k}u^{k-j}(\partial^2 u) \qquad (\partial v)(\partial v)(\partial v)\tilde{v}$$

$$(k-j)! \qquad (k-j)! \qquad (4.9)$$

$$+6\frac{(k-3)!}{u^{k-j}}(\partial u) \qquad (\partial^{2}v)(\partial v)(\partial v)\tilde{v} \qquad (4.8)$$

$$+6\frac{(k-3)!}{(k-j)!} \stackrel{0}{\circ} \times 0 \qquad x^{i} \times l \times m^{i}, l, m}{i^{j-i}l < m}$$

$$+6\frac{(k-3)!}{u^{k-j}}(\partial^{2}u) \qquad (\partial v)(\partial v)(\partial v)\tilde{v} \qquad (4.9)$$

$$(k-j)! \stackrel{0}{\circ} \times 0 \qquad x^{i} \times l \times m^{i}, l, m$$

$$i^{j-1}l = 1 m = 1$$

$$i^{j-1}l = 1 m = 1$$

$$i^{j-1}l < m$$

$$+6\frac{(k-3)!}{(k-3)!} u^{k+1-j} \stackrel{\sum \sum \sum \sum j}{\sum j} \stackrel{\sum j}{\sum j} \stackrel{\sum$$

where we are using the notation  $\tilde{v}_{i_1,...,i_p}$  to express the product of all the vectors  $v_1,\ldots,v_j$ , except  $v_{i_1},\ldots,v_{i_p},$  for  $1\leq i_1,\ldots,i_p\leq j$  distinct from each other and  $1\leq p\leq j$  – 1, and  $\tilde{v}_{i_1,\ldots,i_f}=1$ . Also, the term (4.3) appears only for j = 1, ..., k - 3, the terms (4.4) and (4.5) appear only for  $j=1,\ldots,k-2$ , the terms (4.6) and (4.7) appear only for  $j=2,\ldots,k-1$  the terms (4.8) and (4.9) appear only for j = 3, ..., k and finally, the term (4.10) appears only for j = 4, ..., k + 1.

Without loss of generality, for this term, we will assume that  $4 \le j \le k - 3$  since in this case all terms in formulas (4.3) - (4.10) make sense and for  $j \in \{0, 1, 2, 3\}$  we can do the computation separately and we obtain an estimate compatible with the case  $4 \le j \le k - 3$ . By using formula (2.3) in Lemma 3 and triangle inequality, for  $v_1,\ldots,v_j\in G^{\delta,s+2}$  we have

In Lemma 3 and triangle inequality, for 
$$v_1, ..., v_j \in G^*$$
 we have
$$\|D_j F_1(u_j)(v_j, ..., v_j)\|_{b_j, s+2} \leq 3|a|(k-2) \frac{(k-3)!}{(k-3-j)!} \|u^{k-3-j}(\partial_j u_j)^3(\partial_j^2 u_j)v_j \cdots v_j\|_{b_j, s}$$

$$+ \frac{(k-3)!}{(k-2-j)!} \|u^{k-2-j}(\partial_j u_j)^3 \sum_{j=1}^{2^{-j}} (\partial_j^2 v_j)\tilde{v}\|_{b_j, s}$$

$$+ 3 \frac{(k-3)!}{(k-2-j)!} \|u^{k-2-j}(\partial_j u_j)^2(\partial_j^2 v_j) \otimes (\partial_j v_j)\tilde{v}\|_{b_j, s}$$

$$+ 3 \frac{(k-3)!}{(k-1-j)!} \|u^{k-1-j}(\partial_j u_j)^2(\partial_j^2 v_j) \otimes (\partial_j v_j)\tilde{v}\|_{b_j, s}$$

$$+ 6 \frac{(k-3)!}{(k-1-j)!} \|u^{k-1-j}(\partial_j u_j)(\partial_j^2 u_j) \otimes \sum_{j=1}^{2^{-j}} \sum_{j=1}^{2^{-j}} (\partial_j v_j)(\partial_j v_j)\tilde{v}\|_{b_j, s}$$

$$+ 6 \frac{(k-3)!}{(k-j)!} \|u^{k-j}(\partial_j u_j) \otimes \sum_{j=1}^{2^{-j}} \sum_{j=1}^{2^{-j}$$

Now, from the algebra property in Lemma 4 and (2.2) in Lemma 3, we obtain

Finally, by using lemmas 3 and 4, for  $0 < \delta^i < \delta \le 1$  and  $v_1, \dots, v_i \in G^{\delta, s+2}$ , with s > 1/2, we can estimate

$$\begin{split} &\|D^{j}F_{1}(u_{0})(v_{1},\ldots,v_{j})\|_{\delta',s+2} \\ &\leq \frac{3|a|c^{k}(k-2)e^{-1}}{\delta-\delta^{l}} \|u_{0}\|_{1,s+2} \|v_{1}\|_{\delta,s+2} \cdots \|v_{j}\|_{\delta,s+2} \frac{(k-3)!}{(k-3-j)!} + \\ &+ \frac{4(k-3)!j!}{(j-1)!(k-2-j)!} + \frac{9(k-3)!j!}{(j-2)!(k-1-j)!} + \frac{12(k-3)!j!}{(j-3)!(k-j)!} + \frac{6(k-3)!j!}{(j-4)!(k+1-j)!} \end{split}.$$

By using the fact that  $(k-2) \le 2^{k-3}$  for  $k \in \{3, 4, ....\}$  we obtain

$$\begin{split} &\|D^{j}F_{1}(u_{0})(v_{1},\ldots,v_{j})\|_{\delta',s+2}\\ &\leq \frac{36|a|c^{k}2^{k-3}e^{-1}j!}{\delta-\delta^{j}}\|u_{0}\|_{1,s+2}\|v_{1}\|_{\delta,s+2}\cdots\|v_{j}\|_{\delta,s+2} & k-3\\ &+ k-3 + k-3 + k-3\\ &j-2 + j-3 + j-4 \end{split}$$

Since

$$\|D^{j}F(u_{0})(v_{1},...,v_{j})\|_{\delta',s+2} \leq \frac{AC_{0}^{j}j!}{\delta-\delta^{j}}\|v_{1}\|_{\delta,s+2}\|v_{2}\|_{\delta,s+2}\cdots\|v_{j}\|_{\delta,s+2}, \tag{4.12}$$

where 
$$A = (2|a| + |b+1| + 2|c| + |9a+b| + 3)c^k 2^{3k+2}e^{-1}||u_0||_{1,s+2}^{k+1}$$

Therefore, by Theorem 2 we conclude that the Cauchy problem for the k-abc-equation (1.1) with initial condition  $u(x, 0) = u_0(x)$  has a unique solution, which for  $0 < \delta < 1$  is a holomorphic function in the disc  $D(0,T(1-\delta))$  valued in  $G^{\delta,s+2}$ . Moreover, the lifespan T is given by

$$T = \frac{1}{2e^2AC_0} = \frac{1}{c\|u\|_{1,s+2}^k}, \text{ where } c = e(2|a| + |b+1| + 2|c| + |9a+b| + 3)c^k 2^{3k+3}.$$
 (4.13)

The proof of Theorem 1 is now complete. Q

### The Power Series Method for Nonlocal and Nonlinear ...

#### REFERENCES

- M. S. Baouendi and C. Goulaouic, Remarks on the abstract form of nonlinear Cauchy-Kovalevsky theorems, Comm. in Partial Differential Equations, 2(11), (1977), 1151–1162.
- S. Baouendi and C. Goulaouic, Sharp estimates for analytic pseudodifferential operators and application to Cauchy problems, J. Differential Eqns 48 (1983).
- c. R. Barostichi, A. Himonas and G. Petronilho, A Cauchy-Kovalevsky theorem for nonlinear and nonlocal equations. PROMS Springer-Verlag series, Analysis and Geometry, Volume 127, (2015), 59–68.
- d. R. Barostichi, A. Himonas and G. Petronilho, Autonomous Ovsyannikov theorem and applications to nonlocal evolution equations and systems. J. Funct. Anal. 270 (2016), 330-358.
- R. Beals, D. Sattinger and J. Szmigielski, Multipeakons and the classical moment problem. Adv. Math 154 (2000), 229-25
- f. J. Bourgain, Fourier transform restriction phenomena for certain lattice subsets and applications to non-linear evolution equations. Part II: The KdV equation. Geom. Funct. Anal. 3 (1993), 209-262.
  - g. R. Camassa and D. Holm, An integrable shallow water equation with peaked solitons. Phys. Rev. Lett. 71 i. (1993), 1661-1664.
  - C. Cao, D. Holm and E. Titi, Traveling wave solutions for a class of one-dimensional nonlinear shallow water wave models. J. Dyn. Diff. Eqs 16 (2004), 167-178.
  - G. Coclite and K. Karlsen, On the well-posedness of the Degasperis-Procesi equation. J. Funct. Anal. 233
     i. (2006), 60-91.
  - J. Colliander, M. Keel, G. Staffilani, H. Takaoka and T. Tao, Sharp global well posedness for KdV and modified KdV on R and T. J. AMS 16 (2003), 705-749.
  - Constantin and D. Lannes, The hydrodynamical relevance of the Camassa-Holm and Degasperis-Procesi equations. Arch. Ration. Mech. Anal. 192 (2009), 165-186.
  - Constantin and H. McKean, A shallow water equation on the circle. Comm. Pure Appl. Math. 5 (1999), 949-982
- m. Constantin and W. Strauss, Stability of peakons. Comm. Pure Appl. Math. 53 (2000), 603-610.
- Degasperis and M. Procesi, Asymptotic integrability, in "Symmetry and perturbation theory", editedby A. Degasperis and G. Gaeta, World Sci. Publ., 1999.
- J. Escher and Z. Yin, Well-posedness, blow-up phenomena, and global solutions for the b-equation. J. Reine Angew. Math. 624 (2008), 51-80.
- p. A.S. Fokas, On a class of physically important integrable equations. Phys. D 87 (1995), 145-150.
- q. Fuchssteiner and A.S. Fokas, Symplectic structures, their Bäklund transformations and hereditary symmetries. Phys. D 4 (1981), 47-66.
- Fuchssteiner, Some tricks from the symmetry toolbox for nonlinear equations: generalizations of the Camassa-Holm equation. Phys. D 95 (1996), 229-243.
- K. Grayshan and A. Himonas, Equations with peakon traveling wave solutions. Adv. Dyn. Syst. Appl. 8 (2013), 217-232.
- t. G. Gui, Y. Liu, P. Olver and C. Qu, Wave-breaking and peakons for a modified Camassa-Holm equation.